

Problems on General Relativity: 12* optional

January 9, 2022

Problem 1. Consider the extremal case of the Kottler spacetime ("Schwarzschild–de Sitter") as the value of ΛM^2 reaches the upper limit, that is

$$g^K = -f(r)dt^2 + \frac{dr^2}{f(r)} + r^2(d\theta^2 + \sin^2\theta d\phi^2) \quad (1)$$

$$f(r) = 1 - \frac{2M}{r} - \frac{\Lambda r^2}{3}, \quad (2)$$

$$r \geq 0, \quad \Lambda \geq 0, \quad 9\Lambda M^2 = 1.$$

1. Show, that the function f can be written in the following way

$$f(r) = -\frac{\Lambda}{3r}(r - r_e)^2(r + 2r_e) \quad (3)$$

$$r_e := \frac{1}{\sqrt{\Lambda}}. \quad (4)$$

2. Write the metric tensor in a new, ϵ dependent coordinate system $(\tilde{t}, \tilde{r}, \tilde{\theta}, \tilde{\phi})$, defined as follows

$$t =: \frac{\tilde{t}}{\epsilon}, \quad r =: r_e + \epsilon\tilde{r}, \quad \tilde{\theta} = \theta, \quad \tilde{\phi} = \phi, \quad (5)$$

that is

$$g^K = g_{\tilde{t}\tilde{t}}d\tilde{t}^2 + g_{\tilde{r}\tilde{r}}d\tilde{r}^2 + g_{\tilde{\theta}\tilde{\theta}}d\tilde{\theta}^2 + g_{\tilde{\phi}\tilde{\phi}}d\tilde{\phi}^2. \quad (6)$$

3. Show, that there exist all the limits

$$g_{\tilde{x}\tilde{x}}^{(0)} := \lim_{\epsilon \rightarrow 0} g_{\tilde{x}\tilde{x}}, \quad x = \tilde{t}, \tilde{r}, \tilde{\theta}, \tilde{\phi}. \quad (7)$$

4. and that they define a metric tensor (Near Horizon Geometry)

$$g^{\text{NHG}} := g_{\tilde{t}\tilde{t}}^{(0)}d\tilde{t}^2 + g_{\tilde{r}\tilde{r}}^{(0)}d\tilde{r}^2 + g_{\tilde{\theta}\tilde{\theta}}^{(0)}d\tilde{\theta}^2 + g_{\tilde{\phi}\tilde{\phi}}^{(0)}d\tilde{\phi}^2 = -\Lambda\tilde{r}^2d\tilde{t}^2 + \frac{d\tilde{r}^2}{\Lambda\tilde{r}^2} + \frac{1}{\Lambda}(d\tilde{\theta}^2 + \sin^2\tilde{\theta}d\tilde{\phi}^2) \quad (8)$$

5. Can you argue, that the metric tensor g^{NHG} is not isometric to g^K , however its Ricci tensor still satisfies the Einstein equation

$$R_{\mu\nu}^{\text{NHG}} = \Lambda g_{\mu\nu}^{\text{NHG}}. \quad (9)$$

Problem 2. Consider a fluid with energy density ρ , pressure p and 4-velocity u . Its stress-energy tensor is

$$T^{\mu\nu} = \rho u^\mu u^\nu + p(u^\mu u^\nu + g^{\mu\nu}), \quad (10)$$

where g is the spacetime metric tensor. Show, that the equation

$$\nabla_\mu T^{\mu\nu} = 0, \quad (11)$$

amounts to the following system of equations

$$\nabla_\mu ((\rho + p)u^\mu) = u^\mu \partial_\mu p. \quad (12)$$

$$(\rho + p)a^\mu = -(u^\mu u^\nu + g^{\mu\nu})\partial_\nu p \quad (13)$$

$$a := \nabla_u u$$

(the equation of continuity and the force acting on a fluid particle, respectively.)

Hint: contract (11) with u_ν and with $(u^\mu u_\nu + g^\mu{}_\nu)$, respectively, and notice, that the latter tensor is the projection onto the space orthogonal to u . Remember, that $u^\mu u_\mu = -1$.

Problem 3. Let M_{matt} be a spacetime endowed with a metric tensor

$$g^{\text{matt}} = -d\tau^2 + a^2(\tau) (d\chi^2 + s^2(\chi) (d\theta^2 + \sin^2 \theta d\phi^2)), \quad (14)$$

where s denotes one of the following three functions

$$s(\chi) = \chi, \quad \text{or} \quad s(\chi) = \sin \chi, \quad \text{or} \quad s(\chi) = \sinh \chi. \quad (15)$$

Consider a surface $S_{\text{matt}} \subset M_{\text{matt}}$ such that

$$\chi = \chi_0 = \text{const}, \quad (16)$$

parametrised by free values (τ, θ, ϕ) .

Let M_{Sch} be a spacetime equipped with the Schwarzschild metric tensor

$$g^{\text{Sch}} = -\left(1 - \frac{2M}{r}\right) dt^2 + \frac{dr^2}{1 - \frac{2M}{r}} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \quad (17)$$

$$(18)$$

Consider a surface $S_{\text{Sch}} \subset M_{\text{Sch}}$

$$S_{\text{Sch}} = \{(t(\tau), r(\tau), \theta, \phi) : \tau, \theta, \phi \text{ are free}\}. \quad (19)$$

1. calculate the metric tensor

$$q^{\text{matt}} = q^{\text{matt}}{}_{ij} dx^i dx^j, \quad (x^1, x^2, x^3) = (\tau, \theta, \phi) \quad (20)$$

induced in S_{matt} and

$$q^{\text{Sch}} = q^{\text{Sch}}{}_{ij} dx^i dx^j, \quad (x^1, x^2, x^3) = (\tau, \theta, \phi) \quad (21)$$

induced in S_{Sch} .

2. Derive the conditions on $t(\tau), r(\tau), a(\tau)$ and χ_0 upon which

$$q^{\text{matt}}{}_{ij} = q^{\text{Sch}}{}_{ij}. \quad (22)$$

3. Show that curves $\tau \mapsto (\tau, \chi_0, \theta_0, \phi_0) \in S_{\text{matt}}$ are geodesic in M_{matt} .

4. Assume, that curves $\tau \mapsto (t(\tau), r(\tau), \theta_0, \phi_0) \in S_{\text{Sch}}$ are geodesic in M_{Sch} . Show that the consistency condition is

$$\left(\frac{da}{d\tau}\right)^2 = \frac{A}{a} + B, \quad A, B = \text{const}, \quad A > 0. \quad (23)$$

5. Think of glueing the region of M_{matt} bounded by S_{matt} into M_{Sch} or vice versa. Consider the cases $\frac{da}{d\tau} > 0$ and $\frac{da}{d\tau} < 0$, respectively.